

Simple **H**armonic **M**otion

Many types of time-dependence are **oscillatory**,

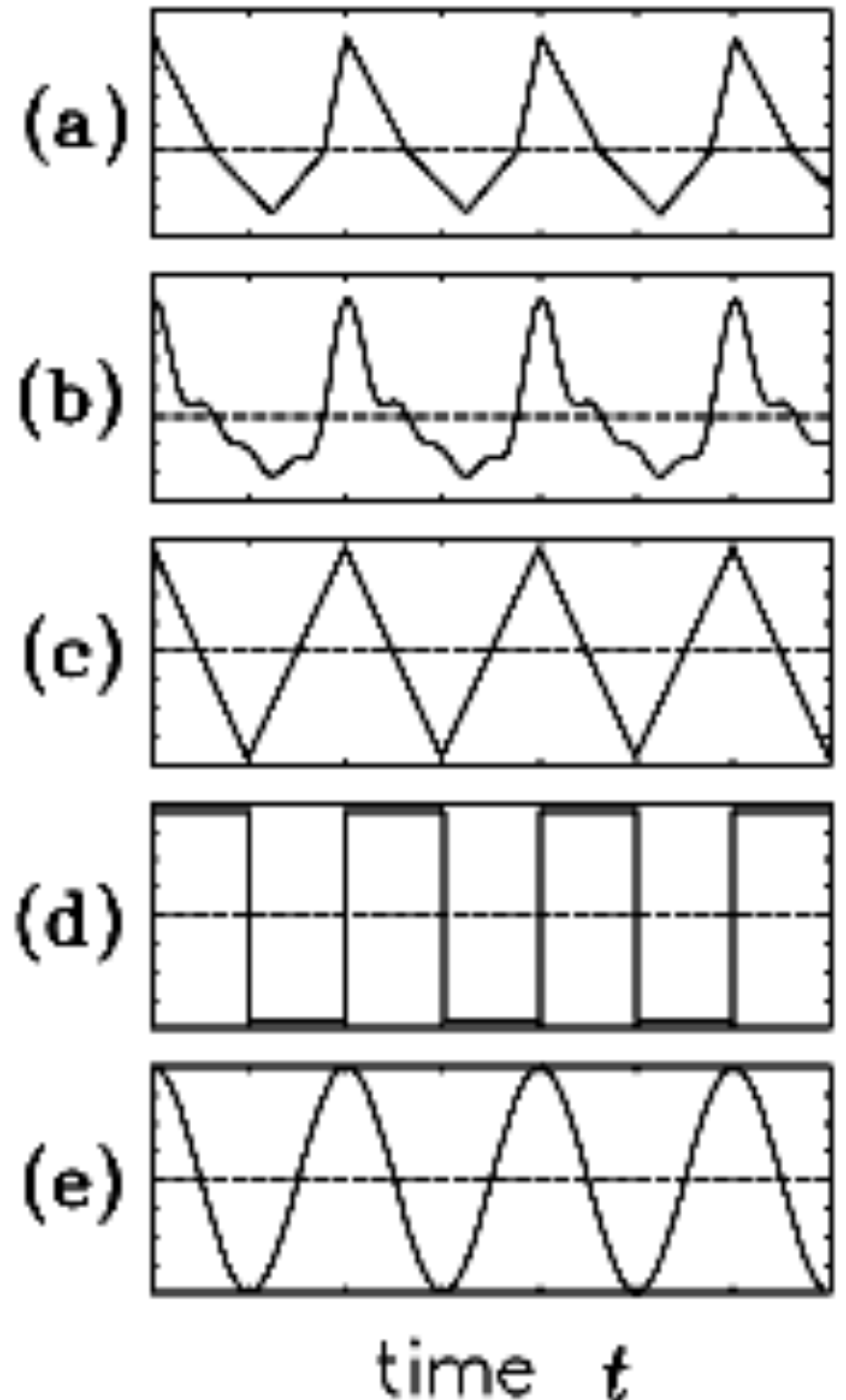
many are **periodic**,

but only one type is “**harmonic**”,

namely,

sinusoidal motion. →

What are its properties?

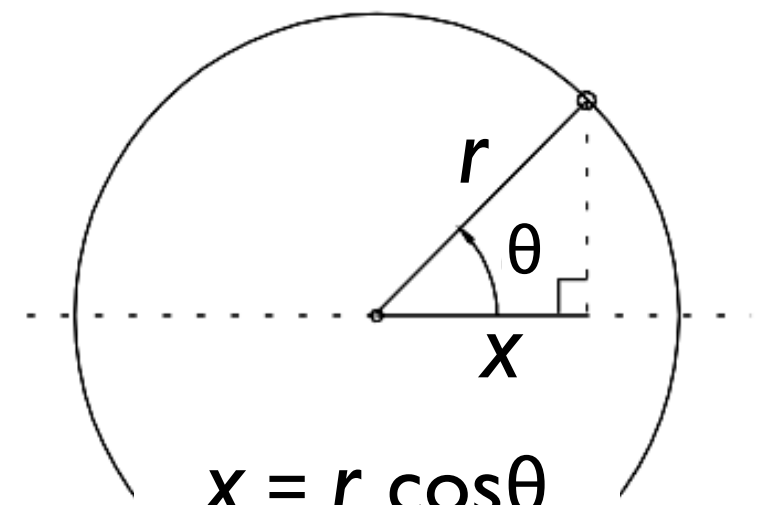
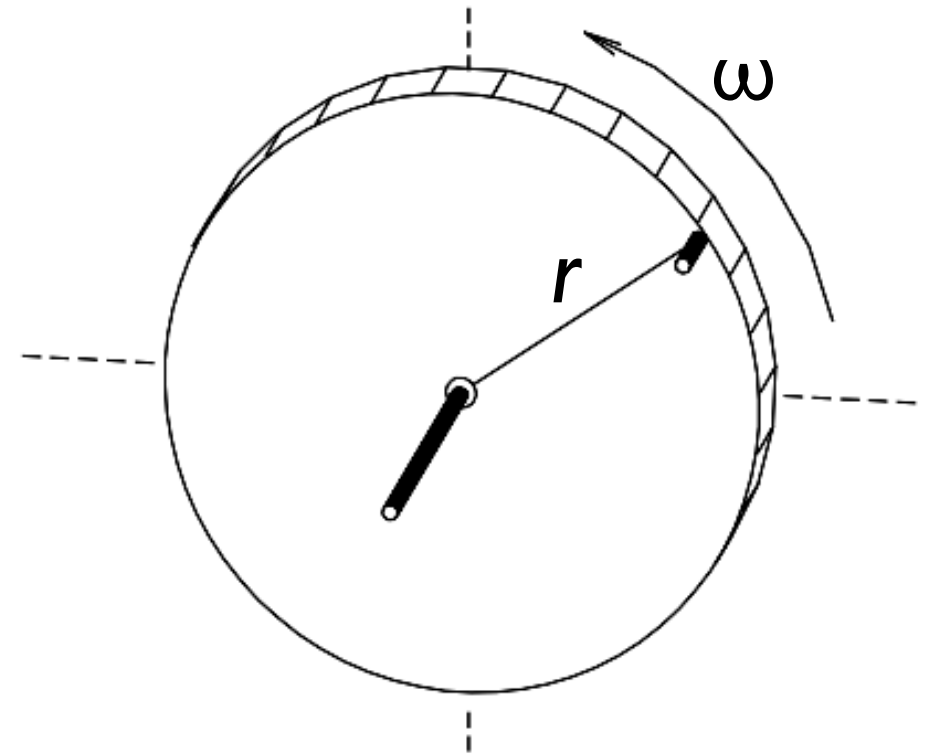


Projecting the Wheel

Picture a wheel spinning at constant angular velocity ω .

Now picture the motion of the **shadow** of a **pin** on the **rim** of the wheel (at high noon on the Equator).

This is called (reasonably) the **projected** motion of the pin.



$$x = r \cos \theta$$

$$\theta = \omega t$$

A Little Mathematics (SWOP):

Small Angles: For $\theta \ll 1$, $\cos(\theta) \approx 1 - \frac{1}{2}\theta^2$
and $\sin(\theta) \approx \theta$.

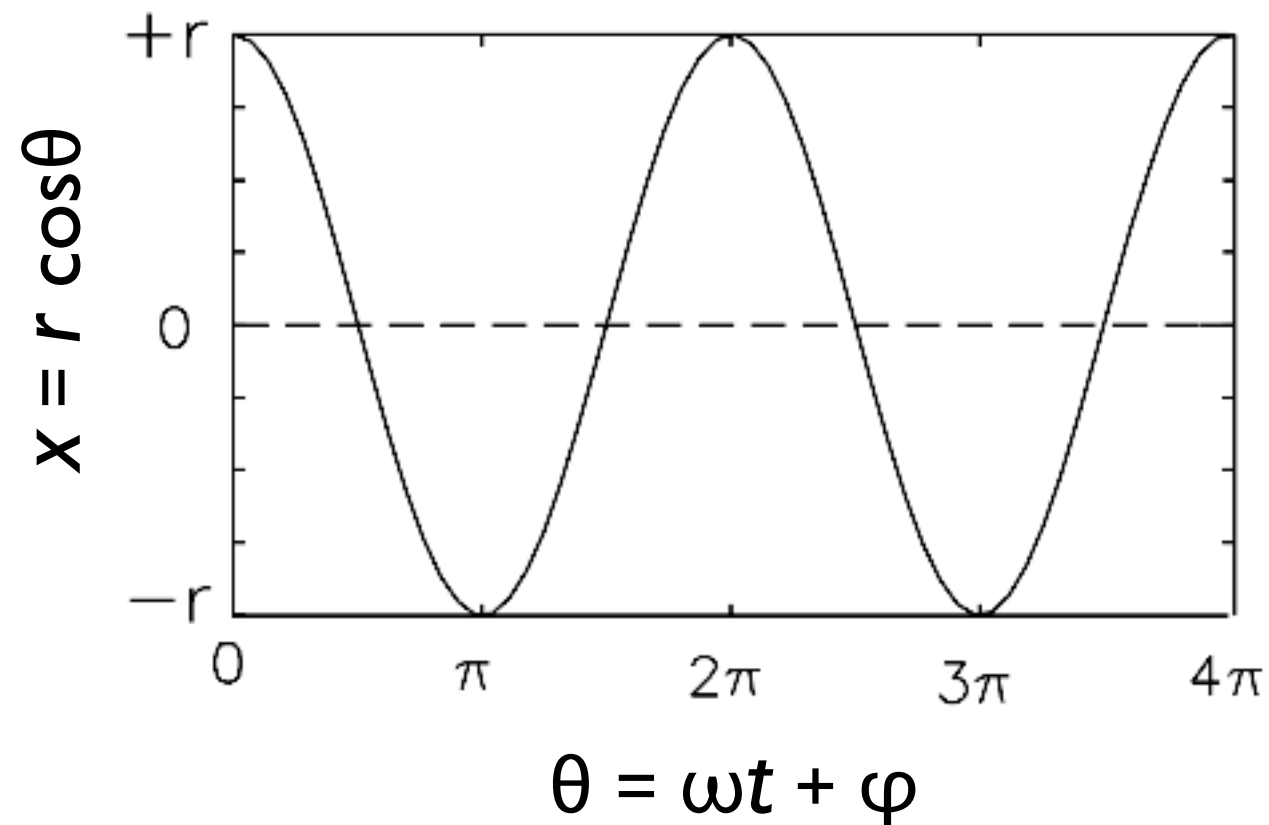
Taylor Series Expansions for Exponential & Sinusoidal Functions:

$$\exp(z) = 1 + z + \frac{1}{2}z^2 + \frac{1}{3!}z^3 + \frac{1}{4!}z^4 + \dots$$

$$\cos(z) = 1 - \frac{1}{2}z^2 + \frac{1}{4!}z^4 - \dots$$

$$\sin(z) = z - \frac{1}{3!}z^3 + \dots$$

Derivatives of the Cosine Function



$$x = r \cos(\omega t + \varphi)$$

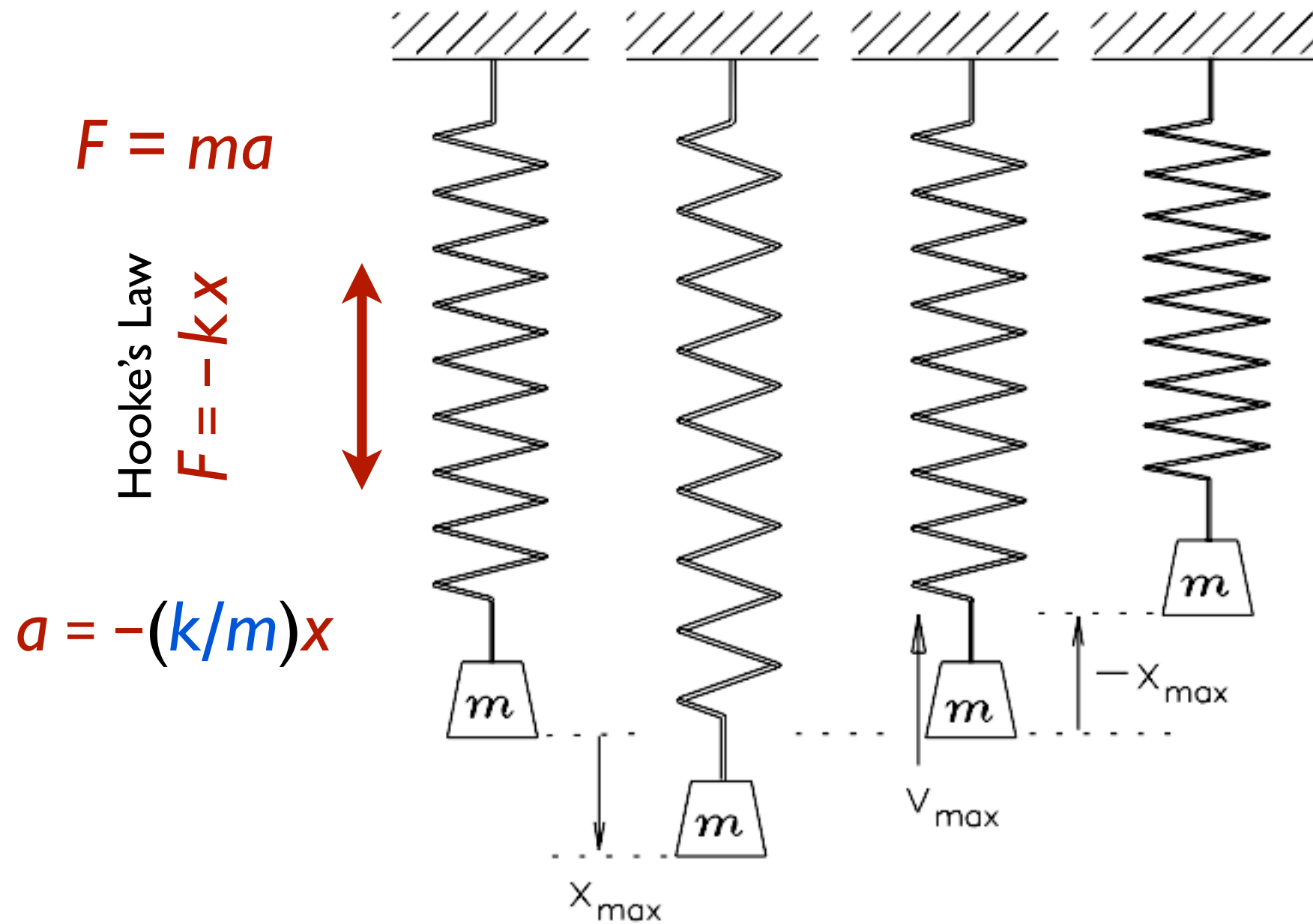
Note:

$$v \equiv dx/dt = -\omega r \sin(\omega t + \varphi)$$

$$d^2x/dt^2 = -\omega^2 x$$

$$a \equiv d^2x/dt^2 = -\omega^2 r \cos(\omega t + \varphi)$$

The Spring Pendulum



So $a \equiv d^2x/dt^2 = -\omega^2x$ if $\omega^2 = k/m$ or

$$\omega = \sqrt{k/m}$$

Our old friend, the Exponential:

Remember $dx/dt = -\kappa x \Leftrightarrow x(t) = x_0 \exp(-\kappa t)$?

The second derivative would be $d^2x/dt^2 = \kappa^2 x$, right?

Well, now we have a new equation $d^2x/dt^2 = -\omega^2 x$, which means the *exponential* function would be a solution if only we could have $\kappa^2 = -\omega^2$.

Of course this is impossible. No real number is negative when squared.

But what if there *were* such a number? Use your *imagination!* $i = \sqrt{-1}$

Then we could solve our differential equation in one step:

$$x(t) = x_0 \exp(i \omega t)$$

— but what does this **mean**?

Complex Exponentials

Taylor Series Expansions for Exponential & Sinusoidal Functions:

$$\exp(z) = 1 + z + \frac{1}{2}z^2 + \frac{1}{3!}z^3 + \frac{1}{4!}z^4 + \dots$$

Recall

$$\cos(z) = 1 - \frac{1}{2}z^2 + \frac{1}{4!}z^4 - \dots$$

$$\sin(z) = z - \frac{1}{3!}z^3 + \dots$$

What if $z = i\theta$?

$$\exp(i\theta) = 1 + i\theta - \frac{1}{2!}\theta^2 - \frac{1}{3!}i\theta^3 + \frac{1}{4!}\theta^4 + \dots$$

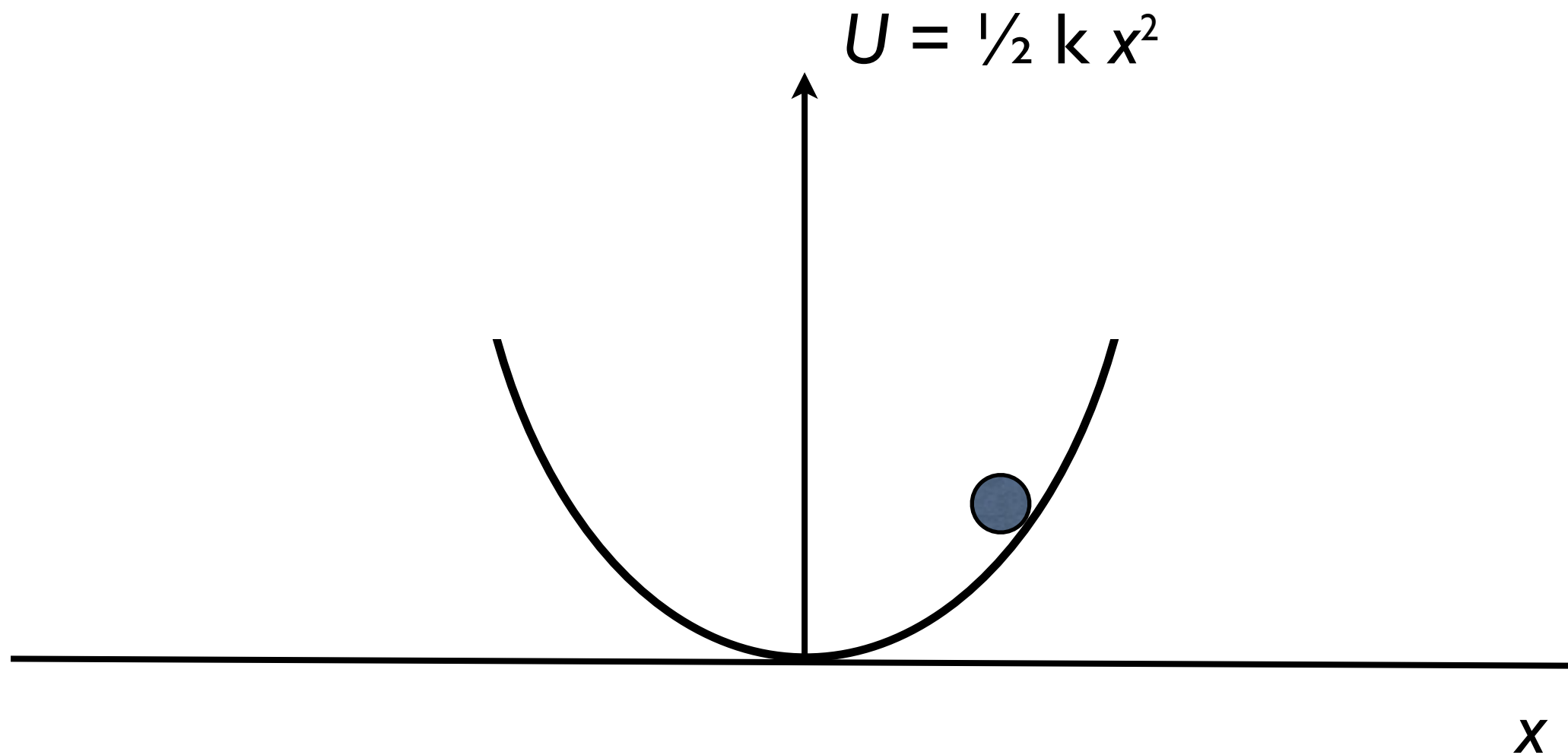
$$\cos(\theta) = 1 - \frac{1}{2!}\theta^2 + \frac{1}{4!}\theta^4 + \dots$$

$$i \sin(\theta) = i\theta - \frac{1}{3!}i\theta^3 + \dots$$

This means

$$e^{i\theta} = \cos \theta + i \sin \theta$$

Quadratic Potential Minimum



$$F = -dU/dx = -kx$$

Simple Harmonic Motion

Linear Restoring Force
(Hooke's Law)

$$F = -kx$$

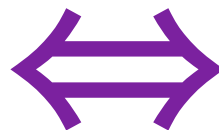


Quadratic Potential Minimum

$$U = \frac{1}{2} k x^2$$

plus

Inertial Factor m



SHM

$$\frac{d^2x}{dt^2} = -\omega^2 x$$



$$x = x_0 \cos(\omega t + \varphi)$$

$$v \equiv \frac{dx}{dt} = -\omega x_0 \sin(\omega t + \varphi)$$

$$a \equiv \frac{d^2x}{dt^2} = -\omega^2 x_0 \cos(\omega t + \varphi)$$

$$\omega^2 = k/m$$

Damped Harmonic Motion:

Viscous damping:

$$d^2x/dt^2 = -\kappa dx/dt \Leftrightarrow v(t) = v_0 \exp(-\kappa t)$$

With a linear restoring force **and** viscous damping, the equation is

$$d^2x/dt^2 = -\kappa dx/dt - \omega^2 x$$

which still might be satisfied by $x(t) = x_0 \exp(Qt)$ with some Q .

Let's try! Plug this $x(t)$ back into the equation, giving

$$Q^2 x = -\kappa Q x - \omega^2 x \quad \text{or} \quad Q^2 + \kappa Q + \omega^2 = 0$$

which has the solution $2Q = -\kappa \pm \sqrt{\kappa^2 - 4\omega^2}$

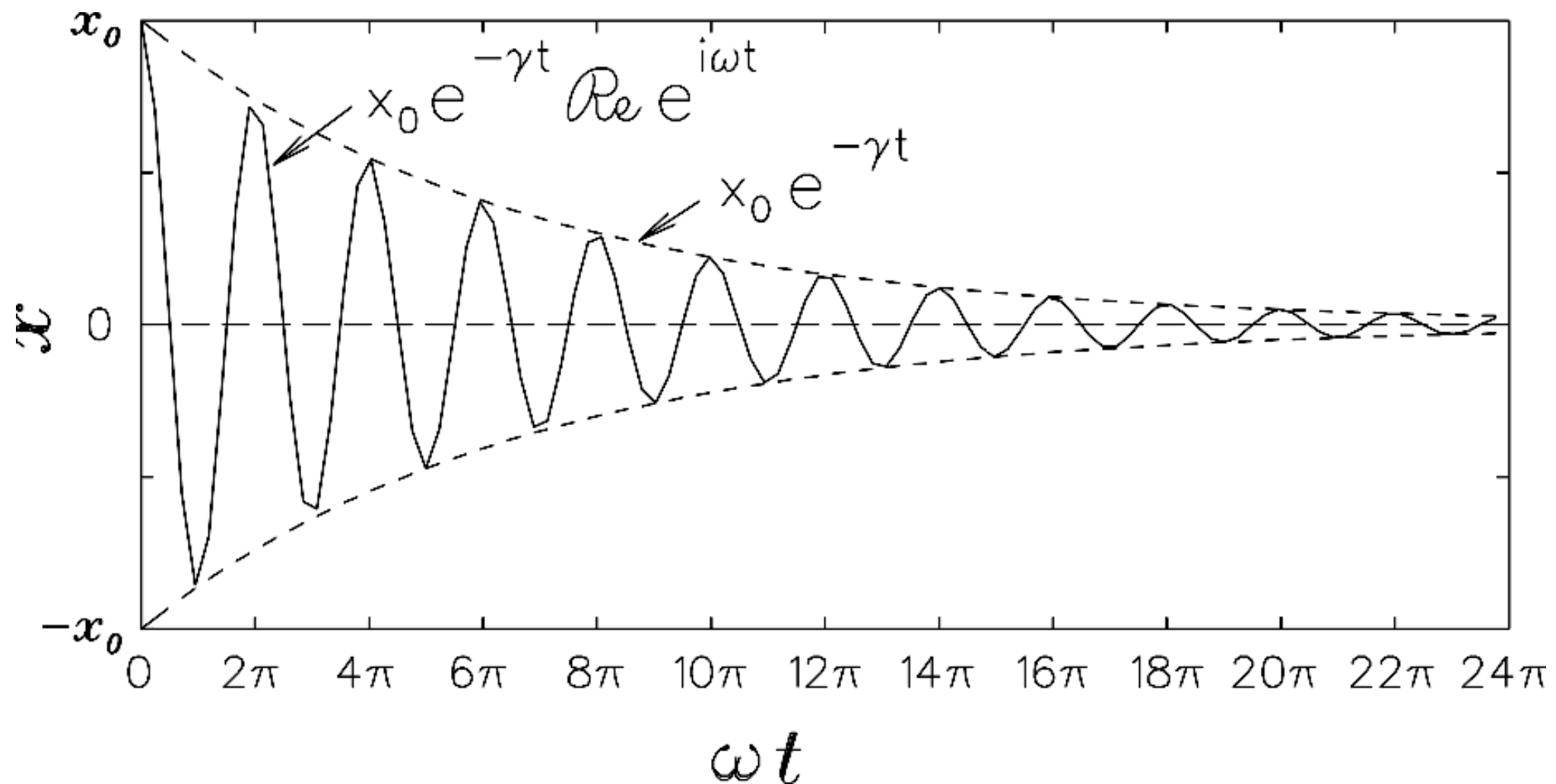
Again, what does this **mean**?

Damped Harmonic Motion:

If $2Q = -\kappa \pm \sqrt{\kappa^2 - 4\omega^2}$ then

$Q = -\frac{1}{2}\kappa \pm i\omega\sqrt{1 - \frac{1}{4}\kappa^2/\omega^2}$, so

$x(t) = x_0 e^{-\frac{1}{2}\kappa t} \exp(\pm i\omega' t)$ where $\omega' = \omega\sqrt{1 - \frac{1}{4}\kappa^2/\omega^2}$



The End

for now